

Search for particle-phonon couplings in ^{65}Cu by the incomplete fusion reaction $^7\text{Li}+^{64}\text{Ni}$

S. Leoni¹, A. Bracco¹, S. Bottoni¹, G. Benzoni¹, F. Crespi¹, L. Pellegrini¹, V. Vandone¹
 N. Mărginean², D. Bucurescu², Gh. Căta-Danil², I. Căta-Danil², D. Deleanu², D. Filipescu²,
 I. Gheorghe², D.G. Ghiță², T. Glodariu², R. Lică², C. Mihai², R.Mărginean²,
 A. Negreț², T. Sava², L. Stroe², S. Toma², R. Șuvăilă², N.V. Zamfir²
 C.A. Ur³
 A. Bruce⁴, O.J. Roberts⁴
 P.H. Regan⁵, Zs. Podolyak⁵, P. J. Mason⁵, C. Townsley⁵

¹ *Università degli Studi di Milano and INFN sez. Milano, Milano, Italy*

² *National Institute for Physics and Nuclear Engineering, Magurele, Romania*

³ *INFN sez. Padova, Padova, Italy*

⁴ *University of Brighton, UK*

⁵ *University of Surrey, UK*

Spokesperson: S. Leoni
Contact Person: N. Mărginean

Abstract

We propose the measurement of excited states in ^{65}Cu populated by the incomplete fusion reaction of ^7Li on ^{64}Ni at energies around the Coulomb barrier. The gamma transitions will be measured using the Bucharest in-beam fast timing array of 8 Ge detectors and 12 LaBr₃(Ce) scintillators, possibly in coincidence with the alpha particles detected with 2 E-DE Si telescopes. The aim of the experiment is two folds: in first place we intend to acquire experience in using incomplete fusion reactions with the weakly bound ^7Li projectile. Such reactions are in fact considered a very powerful spectroscopic tool to get access to highly excited states in n-rich nuclei, at moderately high spin. For this purpose, an excitation function will be performed, by varying the energy of the ^7Li beam between 16 and 22 MeV, in steps of 2 MeV. After determining the optimal beam energy for the population of excited states in ^{65}Cu an in-beam spectroscopy study of this nucleus will be performed. In particular, we intend to focus on particle-phonon coupled states, arising by coupling the unpaired $p_{3/2}$ proton to the 3⁻ octupole phonon of the ^{64}Ni core. By applying fast-timing techniques, the lifetimes of the states of interest will be determined, therefore allowing to estimate their collectivity and to compare with particle-phonon calculations in the weak-coupling limit.

We request 9 days of beam time: 2 for excitation function and 7 for in beam spectroscopy.

Physics motivation

The understanding of particle-phonon and phonon-phonon couplings is a very important issue, since this phenomenon is at the basis of fermionic many-body interacting systems, both in solid state and nuclear physics. In nuclear physics, the coupling between a particle/hole and a vibration is a key ingredient to explain important phenomena, such as the observed reduction of spectroscopic factors, the anharmonicity of vibrational spectra, the damping of Giant Resonances, etc. [1,2].

The best place to search for particle-phonon coupled states is around magic or doubly magic nuclei, where collective vibrations are expected to be quite robust. Experimentally, several indications have been found of discrete states of particle-phonon nature, mostly in medium-heavy nuclei [1], but only in few cases clear evidence has been obtained. In addition, it is still an open question whether states of particle-phonon nature can be considered a general nuclear property, down to the region of medium-light systems with reduced collectivity.

In recent works, evidence has been found for particle/hole-phonon coupled states in $^{47,49}\text{Ca}$, based on the 3^- octupole vibration in ^{48}Ca [3,4]. Furthermore, in ^{67}Cu a fast E3 transition from the $9/2^+$ state to the $3/2^-$ ground state with $B(E3)=17(2)\text{W.u.}$ has been measured in a recent $^{64}\text{Ni}(\alpha,p)^{67}\text{Cu}$ experiment performed in Bucharest [5], suggesting a strong particle-octupole phonon coupling with the 3^- phonon of ^{66}Ni . In all cases, the key information has been the measurement of the lifetime of the states, which has contributed to shed light on the structure of the levels. The n-rich Cu isotopes are particularly interesting in this context, since they provide valuable information on nuclear structure above the $Z=28$ shell closure. In addition, they also present a very peculiar situation, which has been tentatively interpreted as a partial breaking of the particle-phonon coupling model. Figure 1 shows a systematic investigation of the first excited $9/2^+$ levels in the odd $^{59-69}\text{Cu}$ isotopes and their g-ray branching [6]. Dashed lines indicate the energy of the 3^- octupole in the corresponding Ni isotopes. The experimental data on the (a,d) and (^3He ,d) proton stripping reactions shows that in all odd mass Cu isotopes the first $9/2^+$ states around 2.5 MeV have large spectroscopic factors [7,8], consistent with a single-particle character. On the contrary, inelastic scattering of (α,α'), (e,e') and (p,p') [9-13] give, at least in the case of ^{63}Cu and ^{65}Cu , a significantly large E3 strength (~ 20 W.u.), compatible with a coupling with the 3^- phonon of ^{62}Ni and ^{64}Ni , respectively. On the other hand, candidates for $\pi p_{3/2} \otimes \text{Ni}(3^-)$ multiplet were also observed at higher energies, above 3 MeV, in the excitation energy region of the 3^- phonons of Ni. This situation is rather contradictory since it cannot be explained within the usual weak-coupling particle-octupole phonon model. Therefore, further detailed investigation is needed, both theoretically and experimentally. In particular, a firm spin assignment for the states of the multiplet around 3 MeV should be established, together with a more direct determination of the structure of these states (including the $9/2^+$), as follows from lifetime measurements.

In this proposal we intend to focus on ^{65}Cu , which is one proton away from the semi-magic nucleus ^{64}Ni (see Figure 1). In the recent work of Chiara et al. [14], the decay from the $9/2^+$ state has been studied in details by a deep inelastic reaction: four decay branches have been observed, including a very weak, direct decay to the ground state. This $9/2^+$ state has been suggested to arise from a weak coupling between a proton and the ^{64}Ni core. No lifetime measurement has been performed for this state, and no evidence has been found for the states, around 3 MeV, previously interpreted as the $\pi p_{3/2} \otimes \text{Ni}(3^-)$ multiplet.

We therefore propose to further investigate the g-decay of this nucleus, by employing a reaction mechanism that is expected to favor the population of excited states based on collective core excitations, such as incomplete fusion of a ^7Li beam on a ^{64}Ni target, at energies

around the Coulomb barrier. Incomplete fusion reactions are in fact considered a very powerful, little exploited, tool to get access to highly excited states at moderately high spins in n-rich nuclei [15,16].

The aim of the experiment is two folds: first, by performing an excitation function (varying the ${}^7\text{Li}$ beam energy between 16 and 22 MeV), we intend to study the properties of the reactions and to determine the best experimental conditions for the population of excited states in ${}^{65}\text{Cu}$. Then we intend to perform a lifetime analysis of the $9/2^+$ state. If this state has a similar octupole character of the 3^- state in ${}^{64}\text{Ni}$, as suggested by inelastic scattering reactions [10], its lifetime should be of the order of ~ 20 ps (assuming the decay branching reported in Ref. [14]). Such a value can be determined by fast-timing techniques, which are able to provide information in the range of tens of picoseconds to few nanoseconds [17].

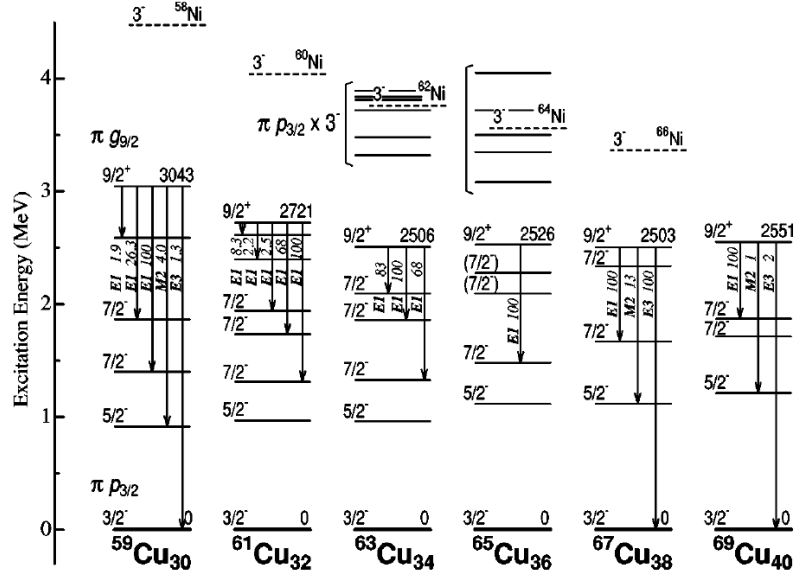


Fig. 1. Energy levels and g -ray branching for the first excited $9/2^+$ states in ${}^{59-69}\text{Cu}$ isotopes. The 3^- octupole states in the corresponding Ni isotopes are indicated by dashed lines. In the case of ${}^{63-65}\text{Cu}$, candidates for $\pi p_{3/2} \otimes \text{Ni}(3^-)$ states are also given [6].

The theoretical interpretation of the experimental results will be done in collaboration with our colleagues Gianluca Colò and Pier Francesco Bortignon of Milano University. It will be based both on a phenomenological approach (originally developed by Bohr and Mottelson [1,3,4]) as well as on a recently developed fully microscopic calculations performed within a self-consistent framework. The latter will be able to provide an exact treatment of the coupling vertex, making use of the whole phonon wave function [18].

It is important to note that the present study forms part of a wider program aiming at a systematic investigation of particle-phonon coupled states in different region of mass and N/Z ratio. It will contribute to extract a precise, quantitative assessment on the coupling strength between particle/hole-states and the low-lying phonon core excitations by comparison with other cases, going from stable to exotic systems. It will also shed light on the observation of an apparent anomalous particle-octupole phonon multiplet, earlier reported in ${}^{65}\text{Cu}$ [13].

Experimental details

We propose to use the incomplete fusion reaction ${}^7\text{Li}$ on ${}^{64}\text{Ni}$ at beam energy of 22 MeV, which is $\sim 30\%$ above the Coulomb barrier. The alpha particles resulting from the incomplete fusion will be detected by 2 Si E-DE telescopes of the ISIS array [19] placed in close geometry around the target to grant 10% detection efficiency. The gamma rays coming in coincidence will be measured using an array of 8 HPGe detectors and 12 LaBr_3 scintillators, with absolute detection efficiency at 1.33 MeV of $\sim 0.8\%$ and $\sim 1\%$, respectively.

A reliable estimate of the $^{64}\text{Ni}(^7\text{Li}, \alpha 2n)^{65}\text{Cu}$ reaction cross section is not a trivial task. However, based on previous experiences [20] and on simple estimates one can expect a cross section $\sigma \sim 50$ mb, mostly peaked at forward angles, around 40° - 50° .

For the excitation function study we plan to employ a ^{64}Ni target of 2 mg/cm^2 on a Au backing of 20 mg/cm^2 , in order to fully stop the recoiling ^{65}Cu isotopes. This will allow to easily identify the γ lines of ^{65}Cu and to study the population of its excited levels. Assuming a ^7Li beam of 3 pA , we expect to measure ~ 15 events/s of α - γ coincidences both with the Ge array and the $\text{LaBr}_3(\text{Ce})$ array. This will be sufficient to perform an excitation function in 2 days, varying the ^7Li beam energy between 16 and 22 MeV, in steps of 2 MeV.

After determining the optimal beam energy for the population of ^{65}Cu , we intend to use a ^{64}Ni target with a thickness of 15 mg/cm^2 , to perform a spectroscopic study of ^{65}Cu . In order to determine the lifetime (or a limit) of the states of interest, triple gamma coincidences will be needed, since one transition observed in the Ge array will be used as a gate (to cleanly select the decay path) and γ - γ coincident transitions detected in the $\text{LaBr}_3(\text{Ce})$ array will be needed to construct the time difference spectrum, according to the method described in Ref. [17]. Assuming a population of 10% for the $9/2^+$ state of interest, a total of ~ 6000 triple gamma coincidences will be collected in 7 days, assuring a meaningful analysis of the time difference spectrum. As a by-product, the use of a thick target will also offer the possibility of performing DSAM measurements for transitions with a lifetime shorter than 1 picosecond.

Our total beam-time request is 9 days

(2 for the excitation function and 7 for in beam measurement)

References

- [1] A. Bohr, B.R. Mottelson, Nuclear Structure, vols. I and II, W.A. Benjamin, 1975.
- [2] P.F. Bortignon, A. Bracco, R.A. Broglia, Giant Resonances: Nuclear Structure at Finite Temperature, Harwood Academic Publishers, New York, 1998.
- [3] D. Montanari et al., Phys. Lett. **B 697**(2011)288.
- [4] D. Montanari et al., submitted to Phys. Rev. C.
- [5] C. Nita et al., to be published.
- [6] M. Asai et al., Phys. Rev. **C62**(2000)054313.
- [7] D. Bucurescu, M. Ivascu, G. Semenescu, and M. Titirici, Nucl. Phys. **A189**(1972)577.
- [8] R. M. Britton and D. L. Watson, Nucl. Phys. **A272**(1976) 91.
- [9] B. G. Harvey et al., Nucl. Phys. **70**(1965)305.
- [10] A. L. McCarthy and G. M. Crawley, Phys. Rev. **150**(1966)935.
- [11] Y. Iwasaki et al., Phys. Rev. C **20**(1979)861.
- [12] A. A. C. Klaasse and V. Paar, Nucl. Phys. **A297**(1978)45.
- [13] A.G. Hartas et al., Nucl. Phys. **A279**(1977)413.
- [14] C.J. Chiara et al., Phys. Rev. C **5**(2012)0234309.
- [15] R.M. Clark et al., Phys. Rev. C **72**(2005)054605.
- [16] G.D. Dracoulis et al., J. Phys. G: Nucl. Part. Phys. **23**(1997)1191.
- [17] J.-M. Régis, G. Pascovici, J. Jolie, M. Rudigier, Nuc. Inst. Meth. **A622**(2010)83.
- [18] G. Colò, H. Sagawa and P.F. Bortignon, Phys. Rev. **C82**(2010)064307.
- [19] E. Farnea et al., Nuc. Inst. Meth. **A400**(1997)87.
- [20] Pfeiffer et al., Nuc. Phys. **A206**(1973)545.